

# Convergence of an axisymmetric finite element $\square$

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## 1) INTRODUCTION

• Let  $\Omega$  be a two-dimensional bounded domain. We suppose that its boundary  $\partial\Omega$  is decomposed into three components  $\Gamma_0$ ,  $\Gamma_D$  and  $\Gamma_N$ :

$$(1.1) \quad \partial\Omega = \overline{\Gamma_0} \cup \overline{\Gamma_D} \cup \overline{\Gamma_N}, \Gamma_0 \cap \Gamma_D = \emptyset, \Gamma_0 \cap \Gamma_N = \emptyset, \Gamma_D \cap \Gamma_N = \emptyset,$$

where  $\Gamma_0$  is the intersection of  $\overline{\Omega}$  with the “axis”  $y = 0$ :

$$(1.2) \quad \Gamma_0 = \overline{\Omega} \cap \{(x, y) \in \mathbb{R}^2, y = 0\}.$$

• Let  $f : \Omega \rightarrow \mathbb{R}$  and  $g : \Gamma_N \rightarrow \mathbb{R}$  be two given functions. We wish to approximate the solution  $u$  of the problem

$$(1.3) \quad -\frac{\partial^2 u}{\partial x^2} - \frac{1}{y} \frac{\partial}{\partial y} \left( y \frac{\partial u}{\partial y} \right) + \frac{u}{y^2} = f \quad \text{in } \Omega$$

$$(1.4) \quad u = 0 \quad \text{on } \Gamma_D$$

$$(1.5) \quad \frac{\partial u}{\partial n} = g \quad \text{on } \Gamma_N$$

where  $n$  is the external normal of the boundary  $\partial\Omega$ .

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$\square$  Presented at the Fourth European Finite Element Fair, Zurich, 2-3 June 2006. Unfinished work. Edition 19 February 2008.

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- The first question is to formulate the problem (1.3)-(1.5) in order to prove the existence and uniqueness. Our variational formulation follows the approach of Mercier and Raugel [MR82] and is briefly recalled in Section 2. By doing this, it is natural to introduce weighted Sobolev spaces  $L_a^2$ ,  $H_a^1$  and  $H_a^2$  associated with axisymmetric problems. The approximation is done with the help of finite elements. We introduce in Section 3 a simplicial conforming mesh  $\mathcal{T}$  composed by vertices (set  $\mathcal{T}^0$ ), edges (set  $\mathcal{T}^1$ ) and triangles (set  $\mathcal{T}^2$ ) and we denote by  $h_{\mathcal{T}}$  the maximal value of the Lebesgue measure of the edges of the mesh  $\mathcal{T}$  :

$$(1.6) \quad h_{\mathcal{T}} = \inf_{a \in \mathcal{T}^1} |a|.$$

We propose a new finite element interpolation based on vertices and defining a discrete space  $H_{\mathcal{T}}^1$  which is “naturally” associated with the Sobolev space  $H_a^1$ . The analysis of this new method is not straightforward. Due to the singular weight  $y$ , it is necessary to use Clément’s interpolate [Cl75] and Section 4 summarizes the essential of what has to be known on this subject. In Section 5, we show that if a function  $u$  belongs to the space  $H_a^2$ , it is possible to define an interpolate  $\Pi_{\mathcal{T}}u$  such that the error  $\|u - \Pi_{\mathcal{T}}u\|$  measured with the norm in space  $H_a^1$ , is of order  $h_{\mathcal{T}}$ . Then the proof of convergence follows classical arguments with Cea’s lemma (see *e.g.* the book [Ci78] of Ciarlet) and is presented in Section 6.

- Some notations:  
 $\text{diam}(K)$  : diameter of the triangle  $K$ .  
where  $|\bullet|$  is the bi-dimensional Lebesgue measure.  
classical Sobolev spaces  
space  $C^0(\bar{\Omega})$ .  
semi-norm in  $H^k(\Lambda)$  Sobolev space:

$$(1.7) \quad |d^k u|^2 \equiv \sum_{\alpha+\beta=k} \left( \frac{\partial^{\alpha+\beta} u}{\partial x^{\alpha} \partial y^{\beta}} \right)^2$$

$$(1.8) \quad |u|_{k, \Lambda}^2 = \int_{\Lambda} |d^k u|^2 dx dy$$

## 2) WEIGHTED SOBOLEV SPACES

- We multiply the equation (1.3) by a test function  $v$  null on the portion  $\Gamma_D$  of the boundary and we integrate by parts relatively to the measure  $y dx dy$ . We introduce by this calculus a bilinear form  $a(\bullet, \bullet)$  and a linear form  $\langle b, \bullet \rangle$  according to

$$(2.1) \quad a(u, v) = \int_{\Omega} y \nabla u \bullet \nabla v \, dx \, dy + \int_{\Omega} \frac{u v}{y} \, dx \, dy$$

$$(2.2) \quad \langle b, v \rangle = \int_{\Omega} f v y \, dx \, dy + \int_{\Gamma_N} g v y \, d\gamma.$$

In consequence of the algebraic expression (2.1) of the bilinear form  $a(\bullet, \bullet)$ , we introduce two notations. If  $u$  is some function  $\Omega \rightarrow \mathbb{R}$ , we define  $u_{\surd}$  and  $u^{\surd}$  as two functions  $\Omega \rightarrow \mathbb{R}$  as

$$(2.3) \quad u_{\surd}(x, y) = \frac{1}{\sqrt{y}} u(x, y), \quad (x, y) \in \Omega$$

$$(2.4) \quad u^{\surd}(x, y) = \sqrt{y} u(x, y), \quad (x, y) \in \Omega.$$

• Following Mercier and Raugel [MR82], we define the three attached Sobolev “axi-spaces”

$$(2.5) \quad L_a^2(\Omega) = \{v : \Omega \rightarrow \mathbb{R}, v^{\surd} \in L^2(\Omega)\}$$

$$(2.6) \quad H_a^1(\Omega) = \{v \in L_a^2(\Omega), v_{\surd} \in L^2(\Omega), (\nabla v)^{\surd} \in (L^2(\Omega))^2\}$$

$$(2.7) \quad H_a^2(\Omega) = \left\{ v \in H_a^1(\Omega), v_{\surd} \surd \surd \surd \in L^2(\Omega), (\nabla v)_{\surd} \in (L^2(\Omega))^2, \right. \\ \left. (d^2 v)^{\surd} \in (L^2(\Omega))^4 \right\}.$$

These spaces are Hilbert spaces associated with the following norms and semi-norms defined according to:

$$(2.8) \quad \|v\|_{0,a}^2 = \int_{\Omega} y |v|^2 \, dx \, dy$$

$$(2.9) \quad |v|_{1,a}^2 = \int_{\Omega} \left( \frac{1}{y} |v|^2 + y |\nabla v|^2 \right) \, dx \, dy$$

$$(2.10) \quad |v|_{2,a}^2 = \int_{\Omega} \left( \frac{1}{y^3} |v|^2 + \frac{1}{y} |\nabla v|^2 + y |d^2 v|^2 \right) \, dx \, dy$$

$$(2.11) \quad \|v\|_{1,a}^2 = \|v\|_{0,a}^2 + |v|_{1,a}^2$$

$$(2.12) \quad \|v\|_{2,a}^2 = \|v\|_{1,a}^2 + |v|_{2,a}^2.$$

We do not need here the expression of the associated scalar products.

• Theorem of trace, hypotheses for  $f$  and  $g$ .

• We observe that the condition

$$(2.13) \quad u = 0 \quad \text{on } \Gamma_0$$

on the axis is completely incorporated inside the choice of the axi-space  $H_a^1(\Omega)$ . We introduce the Sobolev space that takes into account the homogeneous Dirichlet boundary condition (1.4):

$$(2.14) \quad V = \{v \in H_a^1(\Omega), \gamma v = 0 \text{ on } \Gamma_D\}.$$

- Then the problem (1.3)-(1.5) admits the following variational formulation

$$(2.15) \quad \begin{cases} u \in V \\ a(u, v) = \langle b, v \rangle, \forall v \in V. \end{cases}$$

Due to the fact that

$$(2.16) \quad a(v, v) = |v|_{1,a}^2, \quad \forall v \in H_a^1(\Omega),$$

the existence and uniqueness of the solution of problem (2.15) is easy according to the so-called Lax-Milgram-Vishik's lemma and we refer to [MR82] for the study of the ellipticity property.

### 3) A NATURAL AXISYMMETRIC FINITE ELEMENT

- Let  $\mathcal{T}$  be a conforming mesh of the domain  $\Omega$  with triangles. Recall our notations:  $\mathcal{T}^0$  for the set of vertices,  $\mathcal{T}^1$  for edges and  $\mathcal{T}^2$  for triangular elements. We first observe that if we consider a function  $v$  of the form

$$(3.1) \quad v(x, y) = \sqrt{y}(ax + by + c), \quad (x, y) \in K \in \mathcal{T}^2,$$

we have

$$(3.2) \quad \sqrt{y} \nabla v(x, y) = \left( ay, \frac{1}{2}(ax + 3by + c) \right).$$

In other terms, if we denote by  $P_1$  the space of polynomials of total degree less or equal to 1, we have:

$$(3.3) \quad v_{\surd} \in P_1 \implies (\nabla v)_{\surd} \in (P_1)^2.$$

- We denote by  $P_1^{\surd}$  the linear space

$$(3.5) \quad P_1^{\surd} = \{v, v_{\surd} \in P_1\}.$$

We define the degrees of freedom  $\langle \tilde{\delta}_S, v \rangle$  for  $v$  sufficiently regular and  $S$  vertex of the mesh  $\mathcal{T}$  ( $S \in \mathcal{T}^0$ ) by

$$(3.6) \quad \langle \tilde{\delta}_S, v \rangle = v_{\surd}(S), \quad S \in \mathcal{T}^0.$$

We observe that if the vertex  $S$  is not lying on the axis, the number  $\langle \tilde{\delta}_S, v \rangle$  is nothing else that the value  $v(S)$  divided by  $\sqrt{y(S)}$ . If  $S$  is on the axis, consider this point at the origin to fix the ideas and the representation (3.1) joined with (3.6) claims that  $\langle \tilde{\delta}_S, v \rangle = c$ , *id est* is equal to the coefficient of  $\sqrt{y}$  that particularizes the approach. We observe that we still have  $v(S) = 0$  but a non trivial degree of freedom is still present for such a vertex.

**Proposition 1.** Unisolvance property of the axi-finite element.

Let  $K \in \mathcal{T}^2$  be a triangle of the mesh  $\mathcal{T}$ ,  $\Sigma$  the set of linear forms  $\langle \tilde{\delta}_S, \bullet \rangle$  for  $S$  vertex of the triangle  $K$  ( $S \in \mathcal{T}^0 \cap \partial K$ ) and  $P_1^\vee$  defined at relation (3.5). Then the triple  $(K, \Sigma, P_1^\vee)$  that constitutes our axi-finite element is unisolvant.

**Proof of Proposition 1.**

Given three numbers  $\alpha_S \in \mathbb{R}$ , there exists a unique function  $v \in P_1^\vee$  such that

$$(3.7) \quad \langle \tilde{\delta}_S, v \rangle = \alpha_S, \quad S \in \mathcal{T}^0 \cap \partial K.$$

Due to the definition of  $\tilde{\delta}_S$ , the relation (3.7) express that  $v_{\vee}(S) = \alpha_S$  and the hypothesis  $v \in P_1^\vee$  express that  $v_{\vee} \in P_1$ . Then the proof is a consequence of classical arguments for linear finite elements explained *e.g.* in Ciarlet's book.  $\square$

**Proposition 2.** Conformity of the axi-finite element.

The finite element  $(K, \Sigma, P_1^\vee)$  is conforming in space  $\mathcal{C}^0(\overline{\Omega})$ .

**Proof of Proposition 2.**

The property express that given arbitrary values  $\alpha_S \in \mathbb{R}$  for all  $S \in \mathcal{T}^0$ , the function  $v : \Omega \rightarrow \mathbb{R}$  defined by interpolation in each triangle  $K \in \mathcal{T}^2$  by the relation (3.7) is lying in space  $\mathcal{C}^0(\overline{\Omega})$ . The proof is nothing else that the classical  $\mathcal{C}^0$ -conformity of the  $P_1$  finite element:  $v_{\vee} \in P_1$  in each triangle and is defined by its values in each vertex.  $\square$

- We can introduce our discrete space:

$$(3.8) \quad \mathbb{H}_{\mathcal{T}}^\vee = \{v \in \mathcal{C}^0(\overline{\Omega}), v_{\vee}|_K \in P_1, \forall K \in \mathcal{T}^2\}.$$

We have the property:

**Proposition 3.** Conformity in the axi-space  $\mathbb{H}_a^1(\Omega)$ .

The discrete space  $\mathbb{H}_{\mathcal{T}}^\vee$  is included in the axi-space  $\mathbb{H}_a^1(\Omega)$  :

$$(3.9) \quad \mathbb{H}_{\mathcal{T}}^\vee \subset \mathbb{H}_a^1(\Omega).$$

**Proof of Proposition 3.**

It is a direct consequence of the previous property:  $v \in \mathbb{H}_{\mathcal{T}}^\vee$  is continuous then its gradient in the sense of distributions is a classical function. Due to the relation (3.2), this function is clearly in the space  $L^2(\Omega)$ . Of course,  $v_{\vee}$  is continuous then the conditions proposed in (2.6) are all valid.  $\square$

- The discrete space for the approximation of the variational problem (2.15) is simply

$$(3.10) \quad V_{\mathcal{T}} = H_{\mathcal{T}}^{\vee} \cap V.$$

with  $V$  introduced in (2.14). The discrete variational formulation takes the form

$$(3.11) \quad \begin{cases} u_{\mathcal{T}} \in V_{\mathcal{T}} \\ a(u_{\mathcal{T}}, v) = \langle b, v \rangle, \quad \forall v \in V_{\mathcal{T}}. \end{cases}$$

It has a unique solution  $u_{\mathcal{T}} \in V_{\mathcal{T}}$  and the question is now to estimate the error  $\|u - u_{\mathcal{T}}\|$  measured with the norm in the axi-space  $H_a^1(\Omega)$ . For doing this, it is classical to study the interpolation error  $\|u - \Pi_{\mathcal{T}}u\|$  when  $u$  is sufficiently regular and  $\Pi_{\mathcal{T}}u$  is some interpolate of function  $u$ .

#### 4) CLÉMENT'S INTERPOLATION.

- We recall in this section the essential of what to be known about Clément's interpolation [Cl75] in the particular case of affine interpolation with triangles. Let  $\Omega$  be a bounded bidimensional domain as introduced in Section 1. Let  $v$  be a function in space  $L^2(\Omega)$ . Let  $\mathcal{T}$  be a mesh of the domain  $\Omega$  and  $h_{\mathcal{T}}$  introduced in (1.6). We observe also that  $h_{\mathcal{T}}$  is also the maximal diameter of elements in mesh  $\mathcal{T}$  :

$$(4.1) \quad h_{\mathcal{T}} = \sup_{K \in \mathcal{T}^2} \text{diam}(K).$$

Of course, the value  $v(S)$  is not defined for a vertex  $S \in \mathcal{T}^0$  and the interest of Clément's interpolate is to introduce such an approached value even if  $v$  only belongs to the space  $L^2(\Omega)$ .

- First, if  $S \in \Gamma_D$ , we set

$$(4.2) \quad \langle \delta_S^c, v \rangle = 0, \quad S \in \mathcal{T}^0 \cap \Gamma_D.$$

If not, for  $S \in \mathcal{T}^0$ , we introduce the subset  $\Xi_S$  of  $\Omega$  defined by

$$(4.3) \quad \Xi_S = \bigcup_{K \in \mathcal{T}^2, \partial K \supset S} K$$

and presented on Figure 1. The interpolate value  $\langle \delta_S^c, v \rangle$  at the vertex  $S$  is defined by

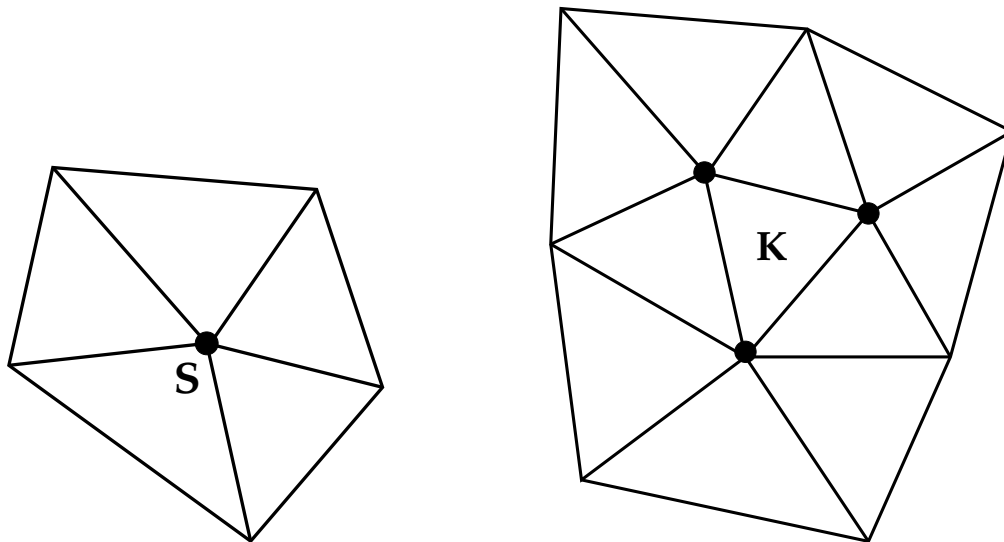
$$(4.4) \quad \langle \delta_S^c, v \rangle = \frac{1}{|\Xi_S|} \int_{\Xi_S} v(x) dx dy, \quad S \in \mathcal{T}^0, S \notin \Gamma_D.$$

- First we introduce the Clement interpolate  $\Pi^c v$  of  $v \in L^2(\Omega)$  with the help of classical  $P_1$  continuous interpolate functions  $\varphi_S$  defined by

$$(4.5) \quad \varphi_S|_K \in P_1, \forall K \in \mathcal{T}^2, \varphi_S(S') = \begin{cases} 1 & \text{if } S' = S \\ 0 & \text{if } S' \neq S. \end{cases}$$

With Clément [Cl75], we set

$$(4.6) \quad \Pi^c v = \sum_{S \in \mathcal{T}^0} \langle \delta_S^c, v \rangle \varphi_S.$$



**Figure 1.** Left: Vicinity  $\Xi_S$  of the vertex  $S \in \mathcal{T}^0$ .

Right: Vicinity  $Z_K$  for a given triangle  $K \in \mathcal{T}^2$ .

- We suppose now that the function  $v$  is a bit more regular. The interest of Clément's interpolation is that all the Ciarlet-Raviart [CR72] classical results for Lagrange interpolation in Sobolev spaces can be extended to Clément's. In order to quantify the result, we suppose in the following that the mesh  $\mathcal{T}$  belongs to a family  $\mathcal{F}$  of meshes such that no infinitesimal angle belongs in the mesh  $\mathcal{T}$ ; in other terms,

$$(4.7) \quad \exists C > 0, \forall \mathcal{T} \in \mathcal{F}, \forall S \in \mathcal{T}^0, \#\{K \in \mathcal{T}^2, K \subset \Xi_S\} \leq C.$$

We introduce also the set  $Z_K$  for a given triangle  $K \in \mathcal{T}^2$  (see again the Figure 1) :

$$(4.8) \quad Z_K = \{L \in \mathcal{T}^2, \bar{K} \cap \bar{L} \neq \emptyset\} = \bigcup_{S \in \mathcal{T}^0, S \subset \partial K} \Xi_S.$$

According to the hypothesis (4.7), we have

$$(4.9) \quad \exists C > 0, \forall \mathcal{T} \in \mathcal{F}, \forall K \in \mathcal{T}^2, \#Z_K \leq C.$$

- Consider now a function  $v \in H^1(Z_K)$ . Then a main results of Clément's contribution can be stated as

$$(4.10) \quad |v - \Pi^c v|_{0,K} \leq C h_{\mathcal{T}} |v|_{1,Z_K}$$

$$(4.11) \quad |v - \Pi^c v|_{1,K} \leq C |v|_{1,Z_K}$$

with a constant  $C > 0$  that does not depend on the particular mesh  $\mathcal{T}$  chosen in the family  $\mathcal{F}$ . If the function  $v$  is more regular ( $v \in H^2(Z_K)$ ), we can consolidate the estimate (4.11):

$$(4.12) \quad |v - \Pi^c v|_{1,K} \leq C h_{\mathcal{T}} |v|_{2,Z_K}.$$

Finally, if  $v$  is globally regular, we have

$$(4.13) \quad \|v - \Pi^c v\|_{0,\Omega} \leq C h_{\mathcal{T}} |v|_{1,\Omega}.$$

## 5) AN INTERPOLATION RESULT

• We suppose in this section that a given function  $u$  belongs to the space  $H_a^2(\Omega)$  defined in (2.7). It is possible to define the value  $u(S)$  for a vertex  $S \in \mathcal{T}^0$  due to the Sobolev embedding Theorem (see *e.g.* Brézis [Br83]) that claims that

$$(5.1) \quad H^2(\Omega) \subset C^0(\overline{\Omega}).$$

The question is now to define or not the number  $\langle \tilde{\delta}_S, u \rangle$  introduced in (3.6).

**Proposition 4.** Lack of regularity.

Let  $u \in H_a^2(\Omega)$  and  $u_{\sqrt{\cdot}}$  introduced in (2.3). Then  $u_{\sqrt{\cdot}}$  belongs to the space  $H^1(\Omega)$  and we have

$$(5.2) \quad \|u_{\sqrt{\cdot}}\|_{1,\Omega} \leq C \|u\|_{2,a}$$

**Proof of Proposition 4.**

We set

$$(5.3) \quad v \equiv u_{\sqrt{\cdot}}$$

and we have the following calculus:

$$(5.4) \quad \nabla v = -\frac{1}{2y\sqrt{y}} u \nabla y + \frac{1}{\sqrt{y}} \nabla u.$$

Then

$$\begin{aligned} \int_{\Omega} |v|^2 \, dx \, dy &\leq \int_{\Omega} \frac{1}{y} |u|^2 \, dx \, dy \leq C \|u\|_{2,a}^2 \\ \int_{\Omega} |\nabla v|^2 \, dx \, dy &\leq 2 \int_{\Omega} \left( \frac{1}{4y^3} |u|^2 + \frac{1}{y} |\nabla u|^2 \right) \, dx \, dy \leq C \|u\|_{2,a}^2. \end{aligned}$$

Due to (2.10), the relation (5.2) is established.  $\square$



- If we derive (formally !) the relation (5.4), we get

$$(5.5) \quad d^2v = \frac{3}{4y^2\sqrt{y}} u \nabla y \cdot \nabla y - \frac{1}{y\sqrt{y}} \nabla u \cdot \nabla y + \frac{1}{\sqrt{y}} d^2u$$

and we have not sufficiently powers of  $y$  to be sure that we obtain a finite result when we integrate the square of  $d^2v$ . In consequence, the function  $v$  is not necessarily continuous. Nevertheless, it is possible to define the Clement interpolate of  $u_{\sqrt{\cdot}}$  relatively to the mesh  $\mathcal{T}$  and due to (5.2), this interpolate has good regularity properties. We define our interpolate  $\Pi u$  by conjugation and we set

$$(5.6) \quad \Pi u = (\Pi^c u_{\sqrt{\cdot}})^{\vee}$$

or equivalently

$$(5.7) \quad \Pi u(x, y) = \sqrt{y} (\Pi^c v)(x, y), \quad (x, y) \in K \in \mathcal{T}^2$$

with  $v(\bullet)$  introduced in (5.3).

- We assume that the mesh  $\mathcal{T}$  admits angles that are aware from 0 and  $\pi$  :

$$(5.8) \quad \begin{cases} \exists(\alpha, \beta), 0 < \alpha < \frac{\pi}{2} < \beta < \pi, \forall \mathcal{T} \in \mathcal{F}, \forall K \in \mathcal{T}^2, \\ \forall \theta \text{ angle in } K, \alpha \leq \theta \leq \beta. \end{cases}$$

We observe that the hypothesis (5.8) clearly implies (4.7). We assume also that the sizes of triangles are quasi-uniform:

$$(5.9) \quad \exists \gamma > 0, \forall \mathcal{T} \in \mathcal{F}, \forall a \in \mathcal{T}^1, \gamma h_{\mathcal{T}} \leq |a| \leq h_{\mathcal{T}}$$

with  $h_{\mathcal{T}}$  introduced at the relation (1.6). We have the following interpolation theorem:

**Theorem 1.** An interpolation result.

We suppose that the mesh  $\mathcal{T}$  belongs to a family  $\mathcal{F}$  that satisfy the above hypotheses (5.8) and (5.9). Let  $u \in H_a^2(\Omega)$  and  $\Pi u$  defined by (5.6). Then we have

$$(5.10) \quad \|u - \Pi u\|_{1,a} \leq C h_{\mathcal{T}} \|u\|_{2,a}.$$

- In order to prepare the technical points of the proof, we introduce, following [MR82] the sub-domains  $\Omega_+$  and  $\Omega_-$  as

$$(5.12) \quad \Omega_+ = \{K \in \mathcal{T}^2, \text{dist}(Z_K, \Gamma_0) > 0\}$$

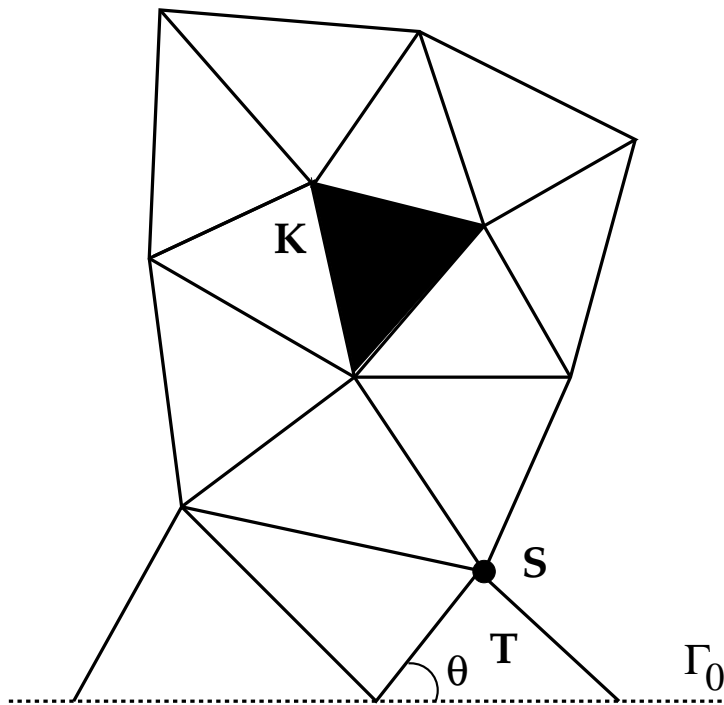
$$(5.13) \quad \Omega_- = \Omega \setminus \Omega_+.$$

**Proposition 5.** Geometrical lemma.

If the family  $\mathcal{F}$  of meshes satisfy the hypotheses (5.8) and (5.9), we have

$$(5.14) \quad \begin{cases} \exists \delta > 0, \forall \mathcal{T} \in \mathcal{F}, \forall K \in \mathcal{T}^2, \\ (K \subset \Omega_+) \implies (\text{dist}(Z_K, \Gamma_0) \geq \delta h_{\mathcal{T}}). \end{cases}$$

$$(5.15) \quad \forall \mathcal{T} \in \mathcal{F}, \forall K \in \mathcal{T}^2, (K \subset \Omega_-) \implies (\forall (x, y) \in K, |y| \leq 2 h_{\mathcal{T}}).$$



**Figure 2.** Triangle element  $K$  that belongs to the sub-domain  $\Omega_+$ .

**Proof of Proposition 5.**

- If  $\text{dist}(Z_K, \Gamma_0) > 0$ , consider a vertex  $S \in \mathcal{T}^0$  such that  $\text{dist}(Z_K, \Gamma_0) = \text{dist}(S, \Gamma_0)$  (see Figure 2). Then the vertex  $S$  belongs to a triangle  $T$  that does **not** belong to the family  $Z_K$  and  $\text{dist}(S, \Gamma_0) \geq \sin \theta \cdot \gamma h_{\mathcal{T}} \geq \gamma \sin \alpha h_{\mathcal{T}}$ . This property establishes the relation (5.14) with  $\delta = \gamma \sin \alpha$ .
- If  $K \subset \Omega_-$  and  $|y| \geq 2 h_{\mathcal{T}}$  for a point  $(x, y) \in K$ , it is clear from the definition of  $Z_K$  and is illustrated by the Figure 2 that the distance between  $Z_K$  and the axis  $\Gamma_0$  is strictly positive, then  $K \subset \Omega_+$  and this contradiction establishes the relation (5.15).  $\square$

**Proof of Theorem 1.**

Our proof is constructed in the same spirit that the pioneering work proposed by Mercier and Raugel.

- We first consider the term of order zero in the error  $\|u - \Pi u\|_{1,a}$  (c.f. the relation (2.11)):

$$\begin{aligned}
 \int_{\Omega} \frac{1}{y} |u - \Pi u|^2 \, dx \, dy &= \int_{\Omega} \frac{1}{y} |u - \sqrt{y} \Pi^{\mathcal{C}} v|^2 \, dx \, dy \\
 &= \int_{\Omega} |v - \Pi^{\mathcal{C}} v|^2 \, dx \, dy = \|v - \Pi^{\mathcal{C}} v\|_{0,\Omega}^2 \\
 &\leq C h_{\mathcal{T}}^2 |v|_{1,\Omega}^2 && \text{due to (4.13)} \\
 &\leq C h_{\mathcal{T}}^2 \|v\|_{2,a}^2 && \text{according to (5.2)}
 \end{aligned}$$

$$(5.16) \quad \int_{\Omega} \frac{1}{y} |u - \Pi u|^2 \, dx \, dy = \|v - \Pi^{\mathcal{C}} v\|_{0,\Omega}^2 \leq C h_{\mathcal{T}}^2 \|u\|_{2,a}^2.$$

• On the other hand, we have

$$\nabla\left(\sqrt{y} (v - \Pi^{\mathcal{C}} v)\right) = \frac{1}{2\sqrt{y}} (v - \Pi^{\mathcal{C}} v) \nabla y + \sqrt{y} \nabla(v - \Pi^{\mathcal{C}} v).$$

Then

$$(5.17) \quad \begin{cases} \int_{\Omega} y |\nabla(u - \Pi u)|^2 \, dx \, dy \leq \\ \leq \int_{\Omega} |v - \Pi^{\mathcal{C}} v|^2 \, dx \, dy + 2 \int_{\Omega} y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy. \end{cases}$$

The first term in the right hand side of (5.17) is majored with the help of estimation (5.16). We focus now on the second term. We have

$$\begin{cases} \int_{\Omega} y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy = \\ = \int_{\Omega_+} y^2 \nabla(|v - \Pi^{\mathcal{C}} v|)^2 \, dx \, dy + \int_{\Omega_-} y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy. \end{cases}$$

• From the relation (5.16), we have for the internal sub-domain  $\Omega_-$  :

$$\begin{aligned}
 \int_{\Omega_-} y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy &\leq 4 h_{\mathcal{T}}^2 \int_{\Omega_-} |\nabla(v - \Pi v)|^2 \, dx \, dy \\
 &\leq C h_{\mathcal{T}}^2 |v|_{1,\Omega}^2 && \text{due to (4.11)} \\
 &\leq C h_{\mathcal{T}}^2 \|u\|_{2,a}^2 && \text{thanks to (5.2)}
 \end{aligned}$$

$$(5.18) \quad \int_{\Omega_-} y |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy \leq C h_{\mathcal{T}}^2 \|u\|_{2,a}^2.$$

• We have in the external part of the domain

$$\int_{\Omega_+} y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy = \sum_{K \in \mathcal{T}^2, K \subset \Omega_+} \int_K y^2 |\nabla(v - \Pi^{\mathcal{C}} v)|^2 \, dx \, dy.$$

We fix  $K \subset \Omega_+$ , we introduce  $(y_{\min}, y_{\max})$  according to

$$y_{\min} = \min_{(x,y) \in K} y, \quad y_{\max} = \max_{(x,y) \in K} y.$$

Then  $y_{\min} \leq y \leq y_{\max}$ ,  $y_{\max} - y_{\min} \leq h_{\mathcal{T}}$  and

$$(5.19) \quad y^2 \leq y_{\max}^2 \leq 2(y_{\min}^2 + h_{\mathcal{T}}^2), \quad (x, y) \in K.$$

Now if  $(x, y) \in Z_K$ , we have, due to the definition of  $Z_K$  and to (5.15):

$$y_{\min} - h_{\mathcal{T}} \leq y, \quad y \geq \delta h_{\mathcal{T}}$$

and

$$(5.20) \quad \frac{y_{\min}^2 + h_{\mathcal{T}}^2}{y^2} \leq \frac{1}{y^2} (2y^2 + 2h_{\mathcal{T}}^2 + h_{\mathcal{T}}^2) \leq 2 + \frac{3}{\delta^2}.$$

• Due to the property (4.12) of Clément's interpolate and to estimate (5.20), we have

$$\begin{aligned} \int_K y^2 |\nabla(v - \Pi^c v)|^2 \, dx \, dy &\leq 2(y_{\min}^2 + h_{\mathcal{T}}^2) \int_K |\nabla(v - \Pi^c v)|^2 \, dx \, dy \\ &\leq 2(y_{\min}^2 + h_{\mathcal{T}}^2) \int_{Z_K} C h_{\mathcal{T}}^2 |d^2 v|^2 \, dx \, dy \\ &\leq C \left(2 + \frac{3}{\delta^2}\right) h_{\mathcal{T}}^2 \int_{Z_K} y^2 |d^2 v|^2 \, dx \, dy \\ &\leq C h_{\mathcal{T}}^2 \int_{Z_K} \left(\frac{1}{y^3} |v|^2 + \frac{1}{y} |\nabla v|^2 + y |d^2 v|^2\right) \, dx \, dy \end{aligned}$$

according to (5.5). Then, taking into account (2.10), we have

$$(5.22) \quad \int_{\Omega_+} y^2 |\nabla(v - \Pi^c v)|^2 \, dx \, dy \leq C h_{\mathcal{T}}^2 \|u\|_{2,a}^2.$$

The inequality (5.11) is a consequence of (5.16), (5.17), (5.18) and (5.22). The theorem is now established.  $\square$

## 6) CONVERGENCE OF THE AXI-FINITE ELEMENT APPROXIMATION.

• We suppose now that the data  $\Omega$ ,  $f$  and  $g$  are chosen in such a way that the solution  $u$  of the variational problem (2.13) belongs to the space  $H_a^2(\Omega)$  :

$$(6.1) \quad u \in H_a^2(\Omega), \quad u \text{ solution of problem (2.13).}$$

Let  $u_{\mathcal{T}} \in H_{\mathcal{T}}^{\vee}$  be the solution of the discrete problem (3.11).

### **Theorem 2.** First order approximation

Under the above hypotheses, we have

$$(6.2) \quad \|u - u_{\mathcal{T}}\|_{1,a} \leq C h_{\mathcal{T}} \|u\|_{2,a}.$$

**Proof of Theorem 2.**

It is a classical consequence of the ellipticity of the functional  $a(\bullet, \bullet)$  and of Céa's lemma. We denote by  $\kappa$  the ellipticity constant of the functional. Thus we get

$$\begin{aligned} \kappa \|u - u_{\mathcal{T}}\|_{1,a}^2 &\leq a(u - u_{\mathcal{T}}, u - u_{\mathcal{T}}) \\ &\leq C a(u - u_{\mathcal{T}}, u - \Pi u) \quad \text{take } v = \Pi u - u_{\mathcal{T}} \in H_{\mathcal{T}}^{\vee} \text{ in (3.11)} \\ &\leq C \|u - u_{\mathcal{T}}\|_{1,a} \|u - \Pi u_{\mathcal{T}}\|_{1,a}. \end{aligned}$$

Then  $\|u - u_{\mathcal{T}}\|_{1,a} \leq C \|u - \Pi u\|_{1,a} \leq C h_{\mathcal{T}} \|u\|_{2,a}$

due to the theorem 1. □

7) REFERENCES.

- [Br85] H. Brézis. *Analyse fonctionnelle. Théorie et applications*, Masson, Paris, 1983.
- [CR72] P.G. Ciarlet, P.A. Raviart. "General Lagrange and Hermite interpolation in  $\mathbb{R}^n$  with applications to finite element methods", *Archive for Rational Mechanics and Analysis*, vol. 46, p. 177-199, 1972.
- [Ci78] P.G. Ciarlet. *The Finite Element Method for Elliptic Problems*, Series "Studies in Mathematics and its Applications", North-Holland, Amsterdam, 1978.
- [Cl75] P. Clément. "Approximation by finite element functions using local regularization", *R.A.I.R.O Analyse numérique*, vol. 9, n°2, p. 77-84, 1975.
- [DD06] F. Dubois, S. Duprey. "Éléments finis naturels pour l'axisymétrie", research report, may 2006.
- [MR82] B. Mercier, G. Raugel. "Résolution d'un problème aux limites dans un ouvert axisymétrique par éléments finis en  $(r, z)$  et séries de Fourier en  $\theta$ ", *R.A.I.R.O Analyse numérique*, vol. 16, n°4, p. 405-461, 1982.