

Construction and analysis of lattice Boltzmann schemes

François Dubois^{*†}

OpenLB Spring School

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* Conservatoire National des Arts et Métiers, LMSSC, Paris

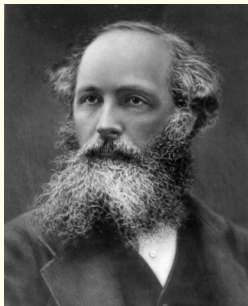
† Laboratoire de Mathématiques d'Orsay, Université Paris-Saclay

velocity distribution at thermodynamic equilibrium

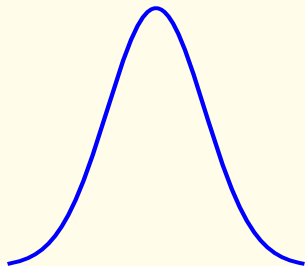
$$f^{\text{eq}}(\mathbf{v}) = \rho \left(\frac{\beta}{2\pi} \right)^{\frac{3}{2}} \exp \left(- \frac{\beta}{2} | \mathbf{v} - \mathbf{u} |^2 \right), \quad \beta = \frac{1}{k_B T}$$

the **temperature T** is a parameter

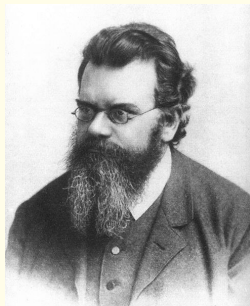
to describe the **variance** of the velocity distribution



James Clerk Maxwell
(1831 - 1879)



Carl Friedrich Gauss
(1777 - 1855)



Ludwig Boltzmann
(1844 - 1906)

Kinetic theory of gases of Maxwell and Boltzmann

$$dm = f(t, x, v) dx dv$$

the gas mass dm is described by a velocity distribution $f(t, x, v)$

Time evolution : Boltzmann equation

$$\frac{\partial f}{\partial t} + v \cdot \nabla_x f = Q(f), \quad x \in \mathbb{R}^3, v \in \mathbb{R}^3, t > 0$$

free advection with velocity v

left hand side $\frac{\partial f}{\partial t} + v \cdot \nabla_x f$

collisions inside the gas

right hand side $Q(f)$

First moments of the distribution f

$$\text{mass } \rho = \int_{\mathbb{R}^3} f(v) dv$$

$$\text{momentum } \rho u = \int_{\mathbb{R}^3} v f(v) dv$$

$$\text{energy } \rho E = \int_{\mathbb{R}^3} \frac{1}{2} |v|^2 f(v) dv$$

$$f(v) = f^{\text{eq}}(v) + \varepsilon f^1(v) + \dots$$

formal derivation of the Navier Stokes equations



Sydney Chapman
(1888 - 1970)

www.npg.org.uk



David Enskog
(1884 - 1947)

www.mech.kth.se

$$\frac{\partial f}{\partial t} + v \cdot \nabla_x f = -\frac{1}{\tau} (f - f^{\text{eq}})$$

collision time $\tau \approx \varepsilon$



Prabhu Lal Bhatnagar
(1912 - 1976)

math.iisc.ernet.in



Eugene Gross
(1926 - 1991)

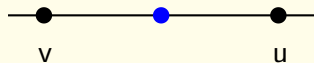
alchetron.com



Max Krook
(1913 - 1985)

[wikipedia](https://en.wikipedia.org/wiki/Max_Krook)

$$\frac{\partial u}{\partial t} + \frac{\partial u}{\partial x} = v^2 - u^2, \quad \frac{\partial v}{\partial t} - \frac{\partial v}{\partial x} = u^2 - v^2$$



Torsten Carleman (1892-1949)

ro.wikipedia.org

Problèmes mathématiques dans la théorie cinétique des gaz

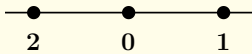
Mittag-Leffler Institute, Stockholm, 1957

$$\frac{\partial f_i}{\partial t} + v_i \cdot \nabla_x f_i = Q_i(f_0, f_1, \dots, f_{q-1}), \quad 0 \leq i < q$$



James Broadwell
(1921-2018)

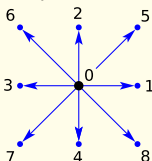
www.nae.edu



D1Q3



Renée Gatignol
(born in 1939)



D2Q9



Henri Cabannes
(1923 - 2016)

ljll.math.upmc.fr


- Boltzmann equation with discrete velocities
- Alternate directions
- Analysis for D1Q3
- Lattice Boltzmann schemes
- Isothermal Navier Stokes
- Thermal Navier Stokes
- Conclusion

ingredients for Multiple Relaxation Times lattice Boltzmann schemes

choice of moments $m = M f$

equilibrium vector function $\Phi(W)$


relaxation matrix S

analysis with a generalization of Chapman Enskog expansion 

inverse problem for Navier Stokes $\Phi(W) = ?$, $S = ?$


isothermal Navier Stokes


D3Q27 has a discrepancy for isothermal Navier Stokes 


D3Q27-2 available for isothermal Navier Stokes 

thermal Navier Stokes

we must impose $\gamma \equiv \frac{c_p}{c_v} = 2$ (2d), $\gamma = \frac{5}{3}$ (3d)

and a Prandtl number satisfying $Pr = 1$ 

D3Q27-2 available for thermal Navier Stokes 

stability has not been studied in this contribution 

numerical experiments are welcomed!

FD, “Equivalent partial differential equations of a lattice Boltzmann scheme”, [Computers and Mathematics with Applications](#), vol. 55, p. 1441-1449, 2008.

FD, “Nonlinear fourth order Taylor expansion of lattice Boltzmann schemes”, [Asymptotic Analysis](#), vol. 127, p. 297-337, 2022.

FD, [zenodo](#) deposit of the “abcd-ns” software, version v0, doi:10.5281/zenodo.6685127, june 2022.

FD and Pierre Lallemand, “On single distribution lattice Boltzmann schemes for the approximation of Navier Stokes equations”, [Communications in Computational Physics](#), vol. 34, p. 613-671, 2023.

FD, Bruce M. Boghosian, Pierre Lallemand, “General fourth-order Chapman-Enskog expansion of lattice Boltzmann schemes”, [Computers & Fluids](#), vol. 266, 106036, 2023.

Thank you for your attention!

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