A new transcendence measure for the values of the exponential function at algebraic arguments

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Abstract

Let $P \in \mathbb{Z}[X] \setminus \{0\}$ be of degree $\delta \geq 1$ and usual height $H \geq 1$, and let $\alpha \in \mathbb{Q}^*$ be of degree $d \geq 2$. Mahler proved in 1931 the following transcendence measure for e^{α} : for any $\varepsilon > 0$, there exists c > 0 such that $|P(e^{\alpha})| > c/H^{\mu(d,\delta)+\varepsilon}$ where the exponent $\mu(d,\delta) = (4d^2 - 2d)\delta + 2d - 1$. Zheng obtained a better result in 1991 with $\mu(d,\delta) = (4d^2 - 2d)\delta - 1$. In this paper, we provide a new explicit exponent $\mu(d,\delta)$ which improves on Zheng's transcendence measure for all $\delta \geq 2$ and all $d \geq 2$. When $\delta = 1$, we recover his bound for all $d \geq 2$, which had in fact already been obtained by Kappe in 1966. Our improvement rests upon the optimization of an accessory parameter in Siegel's classical determinant method applied to Hermite-Padé approximants to powers of the exponential function.

1 Introduction

In this paper, the field of algebraic numbers $\overline{\mathbb{Q}}$ is embedded into \mathbb{C} . Given a number $\xi \in \mathbb{C}$ known to be transcendental over \mathbb{Q} , a classical way to measure how far ξ is from $\overline{\mathbb{Q}}$ is by giving a non-trivial lower bound of $|P(\xi)|$ in terms of the height and degree of $P \in \overline{\mathbb{Q}}[X]$. We then say that we have a *transcendance measure* of ξ . There exist many transcendence measures for classical numbers such as the values of the exponential function at non-zero algebraic numbers, see for instance [3, 10, 14, 16] and the references given in these papers. It is often difficult to compare these results, because sometimes more attention is paid on the height than on the degree, and vice versa, or also because the measure holds when the height or the degree is assumed large enough with respect to the other. In this paper, we shall provide a transcendence measure for e^{α} for all $\alpha \in \overline{\mathbb{Q}}^*$ in which the dependence on the height seems to be the best known so far at this level of generality (Theorem 1 below).

We first introduce some notations. For $\alpha \in \mathbb{K}$ where \mathbb{K} is a number field, we define the house of α as $|\overline{\alpha}| = \max_{\sigma} |\sigma(\alpha)|$ where σ runs through all embeddings of \mathbb{K} into \mathbb{C} ; in other words, $|\overline{\alpha}|$ is the maximum of the moduli of α and of all its Galois conjugates over \mathbb{Q} . Given a polynomial $P(X) = \sum_{j=0}^{d} a_j X^j \in \overline{\mathbb{Q}}[X]$, we set $H(P) := \max_{j=0,\dots,d} |\overline{a_j}|$ its (usual) height. The (usual) height $H(\alpha)$ of $\alpha \in \overline{\mathbb{Q}}$ is defined as H(Q) where $Q \in \mathbb{Z}[X] \setminus \{0\}$ is the minimal polynomial of α over \mathbb{Q} (normalized so that its coefficients are coprime and the leading coefficient is positive). We also let $\deg(\alpha) := \deg(Q)$ and $d(\alpha) \ge 1$ is the denominator of α .

Let p, d, δ be integers such that $d, \delta, p \ge 1$ and $p \ge \delta d$ and let us define

$$\psi(d, \delta, p) := \frac{\delta d^2(p - \delta + 1)}{p - \delta d + 1} + d(p - \delta + 1) - 1.$$

This parameter p in the definition of $\psi(d, \delta, p)$ is accessory but appears naturally in the proof of Theorem 1. It is what enables us to improve on previous bounds. If d = 1, note that $\psi(1, \delta, p) = p$ is independent of δ , but we still require that $p \ge \delta$. For $d \ge 2$, we define $p_1 := \delta d - 1 + \lfloor \delta \sqrt{d^2 - d} \rfloor$ and $p_2 := p_1 + 1$, which are both $\ge \delta d$, and then we set

$$\lambda := \begin{cases} p_1 & \text{if } \psi(d, \delta, p_1) \le \psi(d, \delta, p_2), \\ p_2 & \text{if } \psi(d, \delta, p_2) < \psi(d, \delta, p_1). \end{cases}$$

When d = 1, we set $\lambda := \delta$. We shall prove that λ , which depends on d and δ , minimizes the function $p \mapsto \psi(d, \delta, p)$ amongst all integer values of p such that $p \ge \delta d$ (this is obvious if d = 1 by definition).

The main result of this article is the following new transcendence measure for e^{α} .

Theorem 1. Let \mathbb{K} be a number field, let $\alpha \in \overline{\mathbb{Q}}^*$ be such $[\mathbb{K}(\alpha) : \mathbb{Q}] = d \ge 1$.

For any $\varepsilon > 0$ and any integer $\delta \ge 1$, there exists a constant $c = c(\varepsilon, \alpha, \delta, \mathbb{K}) > 0$ such that for all $H \ge 1$, we have

$$|P(e^{\alpha})| > \frac{c}{H^{\psi(d,\delta,\lambda)+\varepsilon}}$$
(1.1)

for every polynomial $P \in \mathcal{O}_{\mathbb{K}}[X] \setminus \{0\}$ of degree $\leq \delta$ and height $H(P) \leq H$.

We shall also prove the following facts concerning the exponent $\psi(d, \delta, \lambda)$ in Eq. (1.1). For d = 1 and all $\delta \ge 1$, $\psi(1, \delta, \lambda) = \delta$. For all $d \ge 2$ and all $\delta \ge 1$, we have

$$(2d^2 + 2d\sqrt{d^2 - d} - d)\delta - 1 \leq \psi(d, \delta, \lambda) \leq (2d^2 + 2d\sqrt{d^2 - d} - d)\delta - 1 + \frac{d}{\delta\sqrt{d^2 - d} - 1}.$$
 (1.2)

Eq. (1.2) will then be used to show that, for all $d \ge 1$, $\psi(d, \delta, \lambda)$ is an increasing function of $\delta \ge 1$. The constant c can be made completely explicit, and ε can be replaced by an explicit decreasing function of H; the result being unavoidably complicated, we do not state it in this introduction, and we postpone it to Proposition 1 in §3. We simply mention here that the dependence in H is of the form $H^{-\psi(d,\delta,\lambda)-\varpi/\ln\ln(H+2)}$, where $\varpi > 0$ is independent of H; such a dependence already occurs in the celebrated transcendence measure of e obtained by Mahler in [10, p. 135, Satz 3].

We assume in this paragraph that $\mathbb{K} = \mathbb{Q}$ and unless otherwise specified that $\delta \geq 1$ and $d \geq 2$. Another lower bound of the kind in (1.1), *i.e.*, where the attention is paid on H, follows as a special case of the general algebraic independence measure of Lang-Galochkin [6, p. 238] for E-functions: in our situation, it provides the exponent $4\delta d^2$ instead of $\psi(d, \delta, \lambda)$ in (1.1). In [10, pp. 132-133], Mahler had obtained in 1931 a smaller exponent, *i.e.*, $4\delta d^2 - 2\delta d + 2d - 1$ rather than $4\delta d^2$, then improved by Zheng [16] in 1991 with the smaller exponent $(4d^2 - 2d)\delta - 1 =: \mu(d, \delta)$. Notice that when $\delta = 1$, Kappe [9] had already obtained the value $\mu(d, 1)$ in 1966 (¹). Our exponent $\psi(d, \delta, \lambda)$ is always smaller than or equal to all of these exponents because we shall also prove in §2.6 that $\psi(d, \delta, \lambda) \leq (4d^2 - 2d)\delta - 1$ with strict inequality when $\delta \geq 2$ and equality when $\delta = 1$. In fact, for all $\delta \geq 2$ and $d \geq 2$, it can be shown that $\psi(d, \delta, \lambda) \leq (4d^2 - 2d - \frac{1}{4})\delta$; this implies that $\mu(d, \delta) - \psi(d, \delta, \lambda) \geq \frac{\delta}{4} - 1$ which quantifies more precisely the gain we obtain when $\delta \geq 5$. If d = 1, then for all $\delta \geq 1$, the value $\psi(1, \delta, \lambda) = \delta$ is the optimal Dirichlet exponent when $\mathbb{K} = \mathbb{Q}$, and we recover here a well-known result; see [2, Chapter 10].

Since for all fixed $d \ge 1$, $\psi(d, \delta, \lambda)$ is an increasing function of $\delta \ge 1$, it follows that when $\deg(P) \ge 1$ is given, the smallest exponent of H obtainable in (1.1) is for $\delta = \deg(P)$. In particular, when $d \ge 2$:

- (i) When deg(P) = 1, the smallest exponent obtainable in (1.1) is $4d^2 2d 1$, obtained for $\delta = 1$ with $\lambda = 2d 2$ (because $\lfloor \sqrt{d^2 d} \rfloor = d 1$ for all $d \ge 2$).
- (*ii*) When $\deg(P) = 2$, the smallest exponent obtainable in (1.1) is

$$\frac{16d^3 - 16d^2 + d + 1}{2d - 1} = 8d^2 - 4d - \frac{3}{2} + \mathcal{O}\left(\frac{1}{d}\right),$$

obtained for $\delta = 2$ with $\lambda = 4d - 2$ (because $\lfloor 2\sqrt{d^2 - d} \rfloor = 2d - 2$ for all $d \ge 2$).

(*iii*) When $\deg(P) = 3$, the smallest exponent obtainable in (1.1) is

$$\frac{36d^3 - 42d^2 + 7d + 2}{3d - 1} = 12d^2 - 6d - \frac{5}{3} + \mathcal{O}\left(\frac{1}{d}\right)$$

obtained for $\delta = 3$ with $\lambda = 6d - 3$ (because $\lfloor 3\sqrt{d^2 - d} \rfloor = 3d - 2$ for all $d \ge 2$).

(When d is large enough with respect to δ , we have more generally $\lfloor \delta \sqrt{d^2 - d} \rfloor = \delta d - \lceil (\delta + 1)/2 \rceil$.) When $\mathbb{K} = \mathbb{Q}$, (i) coincides with Kappe's result, but (ii) and (iii) are better than any previous known bounds.

¹We warn the reader that when $\mathbb{K} = \mathbb{Q}$ we measure here $|be^{\alpha} - a|$ for $(a, b) \in \mathbb{Z}^2$, not $|e - \frac{a}{b}|$ as Kappe does, hence the -1 in our exponent. Kappe's upper bound also holds for d = 1 (and it is then an equality to 1) and it is still the best known so far for $d \ge 2$ in general. She obtained it by the different but related method of interpolation series, see also [11, Chapter II]. For certain quadratic numbers like $\sqrt{2}$, Kappe's upper bound 11 has been improved to 3 (*i.e.*, $|be^{\sqrt{2}} - a| > c/b^3$) by a different method based on graded Padé approximants to the *E*-function $e^{\sqrt{2}x} + e^{-\sqrt{2}x}$; see [7].

Finally, it seems that the following properties hold, but we did not try to prove them: for all $d \ge 2$, if $\delta \ge 3$ is odd, $\psi(d, \delta, p_1) < \psi(d, \delta, p_2)$, while if $\delta \ge 2$ is even, $\psi(d, \delta, p_2) < \psi(d, \delta, p_1)$.

As already said, Eq. (1.1) improves on Zheng's transcendence measure [16] when $\delta \geq 2$ and $d \geq 2$. Like Zheng and Mahler earlier in [10], our proof uses explicit Hermite-Padé type approximants to powers of the exponential function (which have been known explicitly for a long time) and Siegel's classical "determinant method" to produce linearly independent linear forms in powers of e^{α} . Zheng also considered the parameter p (called m in his paper) which he then took equal to $2\delta d - 1$ in his computations. Our contribution lies in the optimization of p with respect to δ and d, and this leads to our improvement. Another interesting application of Siegel's method is in [13], in which Sorokin computed a very good transcendance measure of π^2 by this method.

Using a different method (*i.e.*, auxiliary functions constructed with Siegel's lemma), Cijsouw [3, Theorem 1] obtained a lower bound $|P(e^{\alpha})| > e^{-c(\alpha)\delta^3}H^{-c(\alpha)\delta^2}$ where $P \in \mathbb{Z}[X] \setminus \{0\}$; his constant $c(\alpha) > 0$ is not explicit but from the proof it seems to be larger than $4d^2$ and also to depend on $H(\alpha)$. Our lower bound is thus again better with respect to the dependence on H but, as we shall infer from the examples deduced from Proposition 1 in §3, his lower bound is better with respect to the dependence on δ . Note that when $H(P) \leq e^{\delta}$, Cijsouw's lower bound has been improved in [14, p. 455, Corollary 3.9].

Besides Zheng's bound, recent works on transcendence measures of e^{α} essentially all focused on the cases where $\alpha \in \mathbb{K}$ and $\mathbb{K} = \mathbb{Q}$ or a quadratic imaginary field, with results similar to those obtained by Mahler for e in [10] or by Baker in [1]. Such results improve on Theorem 1 in these particular cases. We refer for instance to [4, 5, 8] and the references therein.

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2 Proof of Theorem 1 and related results

We first recall basic properties of Hermite-Padé approximants to powers of the exponential. We then prove Theorem 1, and finally we prove various properties stated in the introduction concerning the exponent $\psi(d, \delta, \lambda)$.

2.1 Reminder of Hermite-Padé approximants to power of the exponential

We state here without proofs properties that are immediate applications of the general construction made in [12, pp. 63–69], applied here with m := p + 1 and $\rho_k := (k - 1)\alpha$, $\alpha \in \mathbb{C}^*$. These properties are necessary for the proof of Theorem 1.

For all integers $n \ge 1, p \ge 0$, there exist (explicit) polynomials $\mathcal{P}_{k,\ell} \in \mathbb{C}[x]$ such that

for all $\ell \in \{0, \ldots, p\}$

$$\mathcal{R}_{\ell}(x) := \sum_{k=0}^{p} \mathcal{P}_{k,\ell}(x) e^{k\alpha x}$$

vanishes at order $(p+1)n + \ell$ at x = 0, $\deg(\mathcal{P}_{k,\ell}) = n$ if $k \leq \ell$, $\deg(\mathcal{P}_{k,\ell}) = n - 1$ if $k > \ell$ and

$$\det\left((\mathcal{P}_{k,\ell}(x))_{0\leq k,\ell\leq p}\right) = cx^{(p+1)n}, \quad c\neq 0.$$

For simplicity, we do not write explicitly that \mathcal{R}_{ℓ} and $\mathcal{P}_{k,\ell}$ also depend on α , n and p.

Now, we assume more specifically that $\alpha \in \overline{\mathbb{Q}}^*$, and let \mathbb{K} be a number field over \mathbb{Q} . We define the positive integer $q := \operatorname{lcm}(1, 2, \ldots, p) \times d(1/\alpha)$, so that $q^{(p+1)n+p}n! \mathcal{P}_{k,\ell}(1) \in \mathcal{O}_{\mathbb{Q}(\alpha)} \subset \mathcal{O}_{\mathbb{K}(\alpha)}$. Then, using the explicit formulas and bounds in [12, pp. 66-67], for all $n \geq 1$, all $p \geq 0$, all $\ell = 0, \ldots, p$ and all embeddings σ of $\mathbb{K}(\alpha)$ into \mathbb{C} , we have

$$q^{(p+1)n+p}n! \,\sigma(\mathcal{P}_{k,\ell}(1)) \in \mathcal{O}_{\mathbb{K}(\alpha)},\tag{2.1}$$

$$|q^{(p+1)n+p}n! \sigma(\mathcal{P}_{k,\ell}(1))| \le \left(2q(1+|\sigma(\alpha)|^{-1})\right)^{(p+1)n+p}n!, \tag{2.2}$$

$$|q^{(p+1)n+p}n! \mathcal{R}_{\ell}(1)| \le e^{|\alpha|p(p+1)/2} \cdot \frac{n^{p+1}q^{(p+1)n+p}}{n!^{p}},$$
(2.3)

$$\det\left((P_{k,\ell}(1))_{0\le k,\ell\le p}\right)\neq 0.$$
(2.4)

In (2.2), we can replace $(1 + |\sigma(\alpha)|^{-1})$ by $t_0 := \max_{\sigma} (1 + |\sigma(\alpha)|^{-1}) = 1 + \overline{1/\alpha}$ to obtain an upper bound uniform in σ . Note that $d(1/\sigma(\alpha))$ is independent of σ , which explain the presence of q in (2.1) and (2.2). Note that the upper bound (2.2) can be slightly improved to

$$|q^{(p+1)n+p}n! \sigma(\mathcal{P}_{k,\ell}(1))| \le (2q \max(1, |\sigma(\alpha)|^{-1}))^{(p+1)n+p}(n+1)! \le (2qt)^{(p+1)n+p}(n+1)!$$
(2.5)

by taking directly x = 1 in [12, p. 67, Eq. (120)], and where the above quantity t_0 is now replaced with $t := \max(1, \frac{1}{\alpha})$.

Using [15, p. 80, Lemma 3.11], both quantities $d(1/\alpha)$ and $|1/\alpha|$ can be bounded by $H(\alpha)\sqrt{1 + \deg(\alpha)}$, because $H(\alpha) = H(1/\alpha)$ and $\deg(\alpha) = \deg(1/\alpha)$.

2.2 Proof of Theorem 1

Let $\delta \geq 1$ and $p \geq \delta$. (Later on, we shall even impose that $p \geq d\delta$ to obtain our results.) We set $P(X) := \sum_{k=0}^{\delta} a_k X^k$ with $(a_0, \ldots, a_{\delta}) \in \mathcal{O}_{\mathbb{K}}^{\delta+1} \setminus \{0\}$ such that $H(P) := \max |\overline{a_k}| \leq H$. Let $\alpha \in \overline{\mathbb{Q}}^*$ and define $d := [\mathbb{K}(\alpha) : \mathbb{Q}]$; we have $L_0 := P(e^{\alpha}) \neq 0$ because $e^{\alpha} \notin \overline{\mathbb{Q}}$.

We also set $A_{k,\ell} := q^{(p+1)n+p} n! \mathcal{P}_{k,\ell}(1) \in \mathcal{O}_{\mathbb{K}(\alpha)}$ and $R_{\ell} := q^{(p+1)n+p} n! \mathcal{R}_{\ell}(1)$. We have $\det((A_{k,\ell})_{0 \le k,\ell \le p}) \neq 0$ because $\det((\mathcal{P}_{k,\ell}(1))_{0 \le k,\ell \le p}) \neq 0$.

The $p-\delta+1$ vectors ${}^{t}(a_0,\ldots,a_{\delta},0,\ldots,0), {}^{t}(0,a_0,\ldots,a_{\delta},0,\ldots,0),\ldots,{}^{t}(0,\ldots,0,a_0,\ldots,a_{\delta})$ of \mathbb{C}^{p+1} are \mathbb{C} -linearly independent. We complete them with the δ \mathbb{C} -linearly independent vectors ${}^{t}(A_{0,\ell}, A_{1,\ell}, \ldots, A_{p,\ell})$ $(\ell = 0, \ldots, \delta - 1)$ to form a basis of \mathbb{C}^{p+1} . (2) Consequently, the algebraic integer of $\mathbb{K}(\alpha)$

For every embedding σ of $\mathbb{K}(\alpha)$ into \mathbb{C} , we also have $\sigma(D) \neq 0$ so that

$$\prod_{\sigma} \sigma(D) \in \mathbb{Z} \setminus \{0\}$$

In this product, which is over all such embeddings, we shall distinguish D from the other

 $\sigma(D)$ with $\sigma \neq id$. Let $L_j := \sum_{k=0}^{\delta} a_k e^{(k+j)\alpha}$. On the one hand, by linear combinations of columns, we find

	$ L_0 $	a_1	• • •	a_{δ}	0	• • •	•••	•••	• • •	0
D =	L_1	a_0	•••	$a_{\delta-1}$	a_{δ}	0	•••	•••	•••	0
	:	÷								
	$L_{p-\delta}$	0	•••	0	•••	•••	0	a_0	•••	$\begin{vmatrix} a_{\delta} \\ A_{p,0} \end{vmatrix}$.
	R_0	$A_{1,0}$	• • •	• • •	• • •	• • •	• • •	• • •	• • •	$A_{p,0}$.
	R_1	$A_{1,0} \\ A_{1,1}$	•••	• • •	•••	•••	• • •	•••	• • •	$A_{p,1}$
		÷								:
	$R_{\delta-1}$	$A_{1,\delta-1}$	•••	•••	•••	• • •	• • •	•••	•••	$A_{p,\delta-1}$

Expanding this determinant about its first column and since $L_j = e^{j\alpha}L_0$, the bounds in Eqs. (2.1)-(2.3) imply that

$$|D| \le c_0^{n+1} H^{p-\delta} n!^{\delta} |L_0| + \frac{c_0^{n+1} H^{p-\delta+1}}{n!^{p-\delta+1}}$$
(2.6)

for a constant $c_0 \geq 1$ independent of n and H.

²Strictly speaking, to form a basis of \mathbb{C}^{p+1} , we complete the $p-\delta+1$ first vectors with δ amongst the p+1 vectors ${}^{t}(A_{0,\ell}, A_{1,\ell}, \ldots, A_{p,\ell})$ $(\ell = 0, \ldots, p)$. For ease of writing, we take the δ first ones. In fact, they need not necessarily make up the required basis of \mathbb{C}^{p+1} , but the analysis is completely similar with the δ good vectors. Moreover, this changes nothing to the effective estimates because Eqs. (2.1)-(2.4) are uniform in ℓ .

On the other hand, since

we have

$$|\sigma(D)| \le c_0^{n+1} H^{p-\delta+1} n!^{\delta}$$
(2.7)

where the constant $c_0 \ge 1$ can be taken the same as before (up to increasing it if necessary).

Therefore, from $|D| \prod_{\sigma \neq id} |\sigma(D)| \ge 1$ and with $d := [\mathbb{K}(\alpha) : \mathbb{Q}]$, we deduce that

$$\frac{1}{\left(c_0^{n+1}H^{p-\delta+1}n!^{\delta}\right)^{d-1}} \le |D| \le c_0^{n+1}H^{p-\delta}n!^{\delta}|L_0| + \frac{c_0^{n+1}H^{p-\delta+1}}{n!^{p-\delta+1}}.$$

Hence,

$$|L_0| \ge \frac{1}{c_0^{d(n+1)} H^{(p-\delta+1)(d-1)+p-\delta} n!^{\delta d}} - \frac{H}{n!^{p+1}} =: M$$
(2.8)

The right-hand side satisfies

$$M \ge \frac{1}{2c_0^{d(n+1)}H^{(p-\delta+1)(d-1)+p-\delta}n!^{\delta d}}$$
(2.9)

provided we can choose n (minimal to get a lower bound as large as possible) such that

$$2H^{(p-\delta+1)d} \le n!^{p-d\delta+1} c_0^{-d(n+1)}.$$
(2.10)

Since $H \ge 1$ is arbitrary, a necessary condition for the existence of such an n in all circonstances is that $p - \delta d + 1 > 0$, *i.e.*, that $p \ge \delta d$ because they are integers. We thus now assume that $p \ge \delta d$ and choose n minimal such that (2.10) is satisfied. Combining (2.8) and (2.9) with this value of n, by standard computations (see [12, p. 359] or §3 below), we finally obtain that for all $\varepsilon > 0$, there exists a constant $c = c(\varepsilon, \alpha, \delta, \mathbb{K}) > 0$ independent of H such that

$$|L_0| \ge \frac{c}{H^{\psi(d,\delta,p)+\varepsilon}},$$

where

$$\psi(d, \delta, p) := \frac{\delta d^2(p - \delta + 1)}{p - \delta d + 1} + d(p - \delta + 1) - 1.$$

It remains to find the minimal possible value of $\psi(d, \delta, p)$ under the assumption that $p \ge \delta d$. We recall that when d = 1, $\psi(1, \delta, p) = \delta$ for all $p \ge \delta \ge 1$ so that the minimal value of $\psi(d, \delta, p)$ is achieved at $p = \delta$. We now assume that $d \ge 2$ and $\delta \ge 1$ are fixed. Then the minimum of $x \mapsto \psi(d, \delta, x)$ is attained at

$$x_0 := \delta d - 1 + \delta \sqrt{d^2 - d},$$

and the integers $p_1 := \lfloor x_0 \rfloor = \delta d - 1 + \lfloor \delta \sqrt{d^2 - d} \rfloor$ and $p_2 := \lfloor x_0 \rfloor + 1 = \delta d + \lfloor \delta \sqrt{d^2 - d} \rfloor$ are both admissible to minimize $\psi(d, \delta, x)$ with x an integer (because both are $\geq \delta d$). Therefore defining λ as either p_1 if $\psi(d, \delta, p_1) \leq \psi(d, \delta, p_2)$ or p_2 if $\psi(d, \delta, p_2) < \psi(d, \delta, p_1)$, we obtain that $\psi(d, \delta, p) \geq \psi(d, \delta, \lambda)$ for all $p \geq \delta d$. This completes the proof of Theorem 1.

2.3 The upper bound in (1.2)

Direct computations show that this is true for all $d \ge 2$ if $\delta = 1, 2$, because $\psi(d, 1, \lambda) = 4d^2 - 2d - 1$ and

$$\psi(d,2,\lambda) = \frac{16d^3 - 16d^2 + d + 1}{2d - 1}$$

(because $\lfloor \sqrt{d^2 - d} \rfloor = d - 1$ and $\lfloor 2\sqrt{d^2 - d} \rfloor = 2d - 2$), which are both $\leq (2d^2 + 2d\sqrt{d^2 - d} - d)\delta - 1 + \frac{d}{\delta\sqrt{d^2 - d} - 1}$ for $\delta = 1, 2$ and $d \geq 2$. We now assume that $\delta \geq 3$ and $d \geq 2$.

We start with p_1 . We write

$$p_1 = \delta d + \delta \sqrt{d^2 - d} + x, \quad x \in (-1, 0]$$

and $D := \sqrt{d^2 - d}$ to simplify. Then

$$\psi(d,\delta,p_1) = \frac{\delta d^2(\delta d - \delta + \delta D + x)}{\delta D + x} + d(\delta d + \delta D + x - \delta) - 1$$
$$= 2\delta d^2 + \delta dD - \delta d - 1 + \frac{\delta^2 d^2(d-1)}{\delta D + x} + dx.$$
(2.11)

The function $x \mapsto \frac{\delta^2 d^2(d-1)}{\delta D+x} + dx$ is decreasing on the interval [-1,0] (its derivative is $\frac{dx(2\delta D+x)}{(\delta D+x)^2}$) so that

$$\psi(d, \delta, p_1) \le 2\delta d^2 + \delta d\sqrt{d^2 - d} - \delta d - 1 + \frac{\delta^2 d^2(d-1)}{\delta\sqrt{d^2 - d} - 1} - d$$

Now

$$\frac{\delta^2 d^2 (d-1)}{\delta \sqrt{d^2 - d} - 1} = \delta d \sqrt{d^2 - d} + d + \frac{d}{\delta \sqrt{d^2 - d} - 1}$$

so that

$$\psi(d,\delta,p_1) \le \left(2d^2 + 2d\sqrt{d^2 - d} - d\right)\delta - 1 + \frac{d}{\delta\sqrt{d^2 - d} - 1}$$
(2.12)

For $p_2 = p_1 + 1$, we write

$$p_2 = \delta d + \delta \sqrt{d^2 - d} + x, \quad x \in (0, 1]$$

Then again

$$\psi(d,\delta,p_2) = \frac{\delta d^2(\delta d - \delta + \delta D + x)}{\delta D + x} + d(\delta d + \delta D + x - \delta) - 1$$
$$= 2\delta d^2 + \delta dD - \delta d - 1 + \frac{\delta^2 d^2(d-1)}{\delta D + x} + dx.$$
(2.13)

Since the function $x \mapsto \frac{\delta^2 d^2(d-1)}{\delta D+x} + dx$ is now increasing on the interval [0, 1], we have

$$\psi(d,\delta,p_2) \le 2\delta d^2 + \delta d\sqrt{d^2 - d} - \delta d - 1 + \frac{\delta^2 d^2(d-1)}{\delta\sqrt{d^2 - d} + 1} + d$$

Now

$$\frac{\delta^2 d^2 (d-1)}{\delta \sqrt{d^2 - d} + 1} = \delta d \sqrt{d^2 - d} - d + \frac{d}{\delta \sqrt{d^2 - d} + 1}$$

so that

$$\psi(d,\delta,p_2) \le \left(2d^2 + 2d\sqrt{d^2 - d} - d\right)\delta - 1 + \frac{d}{\delta\sqrt{d^2 - d} + 1}.$$
(2.14)

We remark that the right-hand side of (2.12) is larger than the right-hand side of (2.14), so that

$$\psi(d,\delta,\lambda) \le \left(2d^2 + 2d\sqrt{d^2 - d} - d\right)\delta - 1 + \frac{d}{\delta\sqrt{d^2 - d} - 1}$$

as claimed.

2.4 The lower bound in (1.2)

If $d \geq 2$ and $\delta = 1, 2$, the values $\psi(d, \delta, \lambda)$ recalled in §2.3 (for $\lambda = 2d - 2$ and $\lambda = 4d - 2$ respectively) are both $\geq (2d^2 + 2d\sqrt{d^2 - d} - d)\delta - 1$. For $d \geq 2$ and $\delta \geq 3$, we observe that in (2.11) and (2.13), we obtain lower bounds for $\psi(d, \delta, p_1)$ and $\psi(d, \delta, p_2)$ respectively by taking x = 0. It turns out that these lower bounds are both equal to $(2d^2 + 2d\sqrt{d^2 - d} - d)\delta - 1$.

2.5 $\delta \mapsto \psi(d, \delta, \lambda)$ is increasing

We first recall that λ depends on d and δ , and below we write λ_{δ} for λ . Since $\psi(1, \delta, \lambda_{\delta}) = \delta$, the property holds for d = 1. Let us now assume that $d \geq 2$. We shall prove that for all $\delta \geq 1$, we have $\psi(d, \delta + 1, \lambda_{\delta+1}) > \psi(d, \delta, \lambda_{\delta})$. For this, it is enough to prove that the left-hand side of (1.2) with δ replaced by $\delta + 1$ is larger than the right-hand side of (1.2) with δ . To prove this, we simply observe that

$$(2d^2 + 2d\sqrt{d^2 - d} - d)(\delta + 1) - (2d^2 + 2d\sqrt{d^2 - d} - d)\delta - \frac{d}{\delta\sqrt{d^2 - d} - 1}$$

= $2d^2 + 2d\sqrt{d^2 - d} - d - \frac{d}{\delta\sqrt{d^2 - d} - 1}$
\ge $2d^2 + 2d\sqrt{d^2 - d} - d - \frac{d}{\sqrt{d^2 - d} - 1} > 0$

for all $d \ge 2$ and all $\delta \ge 1$.

2.6 The inequality $\psi(d, \delta, \lambda) \leq 4\delta d^2 - 2\delta d - 1$

If d = 1, $\psi(1, \delta, \lambda) = \delta \leq 2\delta - 1$ for all $\delta \geq 1$. If $d \geq 2$, we simply observe that $\psi(d, \delta, \lambda) \leq \psi(d, \delta, 2\delta d - 1) = 4\delta d^2 - 2\delta d - 1$. The inequality is an equality when $\delta = 1$ and is strict if $\delta \geq 2$ because then $\lambda \leq \delta d + \lfloor \delta \sqrt{d^2 - d} \rfloor < 2\delta d - 1$. Notice that $p := 2\delta d - 1$ is exactly the choice made by Zheng.

3 A completely explicit version of Theorem 1

3.1 The main statement

The goal of this section is to make completely explicit the constant c in Theorem 1. We shall prove the following proposition, which, in fact, subsumes mess Theorem 1.

Proposition 1. Let \mathbb{K} be a number field, let $\alpha \in \overline{\mathbb{Q}}^*$ be such $[\mathbb{K}(\alpha) : \mathbb{Q}] =: d \ge 1$.

For any $H \ge 1$, any integer $\delta \ge 1$, and integer p such that $p \ge d\delta$, and any $P \in \mathcal{O}_{\mathbb{K}}[X] \setminus \{0\}$ of degree $\le \delta$ and height $H(P) \le H$, we have

$$|P(e^{\alpha})| \ge \frac{1}{(2a^d)^{1+\frac{\delta d}{p-d\delta+1}} (b^{d+\frac{\delta d^2}{p-d\delta+1}})^{1+v^2} u^{(d+\frac{\delta d^2}{p-d\delta+1})\frac{4\ln(b)}{\ln\ln(u+2)}} H^{\psi(d,\delta,p)}},$$
(3.1)

where

$$\begin{aligned} q &:= \operatorname{lcm}(1, 2, \dots, p) \times d(1/\alpha), & t &:= \max(1, \overline{|1/\alpha|}) \\ a &:= (p+1)! e^{|\alpha|p(p+1)/2} (2qt)^{p\delta}, & b &:= (2qt)^{(p+1)\delta}, \\ u &:= \left(2a^d H^{(p-\delta+1)d}\right)^{1/(p-d\delta+1)}, & v &:= b^{d/(p-d\delta+1)}e, \\ \psi(d, \delta, p) &:= \frac{\delta d^2(p-\delta+1)}{p-\delta d+1} + d(p-\delta+1) - 1. \end{aligned}$$

After simplifications, the dependence on ${\cal H}$ turns out to be of the standard "Mahlerian" form

$$H^{-\psi(d,\delta,p)-\varpi/\ln\ln(H+2)}$$

for some $\varpi > 0$ independent of H. In [16, Theorem 1], when $\mathbb{K} = \mathbb{Q}$, Zheng has obtained a lower bound of the form $H^{-(4\delta d^2 - 2\delta d - 1) - c\delta^2 / \ln \ln(H+2)}$, where c > 0 is an unspecified constant. We recall that, with respect to the accessory parameter p, the minimal value of $p \mapsto \psi(d, \delta, p)$ is obtained at $p_1 := \delta d - 1 + \lfloor \delta \sqrt{d^2 - d} \rfloor$ or $p_2 = \delta d + \lfloor \delta \sqrt{d^2 - d} \rfloor$.

Proof. The proof follows the same steps as that of Theorem 1, except that we pay attention to effectivity and explicit bounds. For this we shall use the explicit bounds recalled in

Eqs. (2.1)-(2.3) in §2.1. We can then replace (2.6) by the completely explicit bound, valid for all $n \ge 1, p \ge \delta \ge 1, H \ge 1$:

$$|D| \le c_1 c_2^n H^{p-\delta} n!^{\delta} |L_0| + \frac{c_3 c_4^n n^{p+1} H^{p-\delta+1}}{n!^{p-\delta+1}}$$

where

$$c_1 := (p+1)! e^{|\alpha|p} (2qt)^{p\delta}, \quad c_2 := (2qt)^{(p+1)\delta}, c_3 := (p+1)! e^{|\alpha|p(p+1)/2} q^p (2qt)^{p(\delta-1)}, \quad c_4 := q^{p+1} (2qt)^{(p+1)(\delta-1)}.$$

Similarly, we have the explicit bound

$$|\sigma(D)| \le c_5 c_2^n H^{p-\delta+1} n!^{\delta},$$

for all embeddings σ of $\mathbb{K}(\alpha)$ into \mathbb{C} , where $c_5 = (p+1)!(2qt)^{p\delta}$. We now make a few simplifications. We have $n^{p+1} < (2t)^{(p+1)n}$ for all $n \ge 1$ so that $n^{p+1}c_4^n < ((2t)^{p+1}c_4)^n = c_2^n$. Moreover, $\max(c_1, c_3, c_5) < c_6$ where

$$c_6 := (p+1)! e^{|\alpha|p(p+1)/2} (2qt)^{p\delta}$$

Therefore,

$$|D| \le c_6 c_2^n H^{p-\delta} n!^{\delta} \left(|L_0| + \frac{H}{n!^{p+1}} \right)$$
(3.2)

and

$$|\sigma(D)| \le c_6 c_2^n H^{p-\delta+1} n!^{\delta}$$

for all embeddings σ of $\mathbb{K}(\alpha)$ into \mathbb{C} . For simplicity, we now set $b := c_2$ and $a := c_6$.

From the lower bound $|D| \prod_{\sigma \neq id} |\sigma(D)| \ge 1$, where the product is over all such embeddings, we deduce that

$$\frac{1}{\left(ab^{n}H^{p-\delta+1}n!^{\delta}\right)^{d-1}} \leq |D| \leq ab^{n}H^{p-\delta}n!^{\delta}\left(|L_{0}| + \frac{H}{n!^{p+1}}\right)$$

where $d := [\mathbb{K}(\alpha) : \mathbb{Q}]$. Hence,

$$|L_0| \ge \frac{1}{a^{d}b^{dn}H^{(p-\delta+1)(d-1)+p-\delta}n!^{\delta d}} - \frac{H}{n!^{p+1}} =: M$$
(3.3)

The right-hand side of (3.3) satisfies

$$M \ge \frac{1}{2a^{d}b^{dn}H^{(p-\delta+1)(d-1)+p-\delta}n!^{\delta d}}$$
(3.4)

provided we can choose n (as small as possible to get a lower bound as large as possible) such that

$$2a^{d}H^{(p-\delta+1)d} \le n!^{p-d\delta+1}b^{-dn}.$$
(3.5)

Since $H \geq 1$ is arbitrary, a necessary condition for the existence of such an n in all circumstances is that $p - d\delta + 1 > 0$, *i.e.*, that $p \geq d\delta$ because d, δ, p are integers. We thus now assume that $p \geq d\delta$ and ideally we choose the minimal value of n, say n_{\min} , such that (3.5) is satisfied. Finding the exact expression on n_{\min} is difficult but we can find an upper bound for n_{\min} as follows. Since $n! \geq n^n e^{-n}$ for all $n \geq 1$, the minimal value of n (denoted by n_0 from now on) such that

$$2a^{d}H^{(p-\delta+1)d} \le n^{(p-d\delta+1)n}e^{-(p-d\delta+1)n}b^{-dn}$$

is such that $n_{\min} \leq n_0$. Hence (3.5) holds with $n = n_0$, and thus (3.4) as well.

Our task is now to obtain an upper bound for n_0 . By definition of n_0 , we have

$$(n_0 - 1)^{n_0 - 1} < uv^{n_0 - 1}$$

where

$$u := \left(2a^d H^{(p-\delta+1)d}\right)^{1/(p-d\delta+1)} > 1, \quad v := b^{d/(p-d\delta+1)}e > e$$

The integer n_0 is obviously $\geq v + 1$ because $v^v < uv^v$. We then set n = v + x + 1 for some $x \geq 1$ so that for all $s \in (0, 1)$, we have:

$$\frac{(n-1)^{n-1}}{uv^{n-1}} = \frac{(v+x)^{v+x}}{uv^{v+x}}$$

$$\geq \frac{(v+x)^{(v+x)(1-s)}}{u} \quad \text{provided } v+x \geq v^{1/s}$$

$$\geq \left(\frac{x^x}{u^{1/(1-s)}}\right)^{1-s} \quad \text{because } x \geq 1/e$$

$$\geq 1 \quad \text{provided } x \geq 1 + \frac{2\ln(u^{1/(1-s)})}{\ln\ln(u^{1/(1-s)}+2)}.$$

In the last two lines, we use two elementary facts: 1) the function $v \mapsto (v+x)^{v+x}$ increases for $[0, +\infty)$ when $x \ge 1/e$, and 2) for any h > 1, if $x \ge 1 + \frac{2\ln(h)}{\ln\ln(h+2)}$, then $x^x \ge h$. The three assumptions on x are then fulfilled with

$$x_s := v^{1/s} - v + \frac{2\ln(u)}{(1-s)\ln\ln(u^{1/(1-s)} + 2)} + 1.$$

It follows that for all $s \in (0, 1)$, we have

$$n_0 \le v + x_s = v^{1/s} + \frac{2\ln(u)}{(1-s)\ln\ln(u^{1/(1-s)}+2)} + 1.$$

Taking s = 1/2, we deduce the upper bound:

$$n_0 \le v^2 + \frac{4\ln(u)}{\ln\ln(u^2 + 2)} + 1 \le v^2 + \frac{4\ln(u)}{\ln\ln(u + 2)} + 1.$$
(3.6)

Since (3.5) is satisfied with $n = n_0$, we have

$$n_0! > (2a^d H^{(p-\delta+1)d})^{1/(p-d\delta+1)} b^{dn_0/(p-d\delta+1)}.$$

Substituting this into (3.4), we deduce from (3.3) that

$$|L_0| \ge \frac{1}{(2a^d)^{1+\delta d/(p-d\delta+1)} (b^{d+\delta d^2/(p-d\delta+1)})^{n_0} H^{\psi(d,\delta,p)}}$$

Using the upper bound for n_0 in (3.6), we then obtain

$$|L_0| \ge \frac{1}{(2a^d)^{1+\frac{\delta d}{p-d\delta+1}} \left(b^{d+\frac{\delta d^2}{p-d\delta+1}}\right)^{1+\nu^2} u^{(d+\frac{\delta d^2}{p-d\delta+1})\frac{4\ln(b)}{\ln\ln(u+2)}} H^{\psi(d,\delta,p)}},$$

which is the expected lower bound (3.1).

3.2 Some consequences of Proposition 1

In what follows, $\mathbb{K} = \mathbb{Q}$. We present without details three examples of application of Proposition 1.

a) If $\alpha = 1$, we have d = 1, $p = \delta$ and $\psi(d, \delta, p) = \delta$, t = 2, $q = \operatorname{lcm}(1, 2, \dots, \delta) \approx e^{\delta}$, $a \approx e^{\delta^3}$, $b \approx e^{\delta^3}$, $u \approx H^d e^{\delta^3}$, $v \approx e^{d\delta^3}$. We deduce that there exist two absolute constants $c_1 > 0, c_2 > 0$ such that for any $H \ge 1$, any integer $\delta \ge 1$ and any vector $(a_0, \dots, a_{\delta}) \in \mathbb{Z}^{\delta+1} \setminus \{0\}$ with $\max |a_k| \le H$, we have

$$\left|\sum_{k=0}^{\delta} a_k e^k\right| \ge \frac{\exp(-\exp(c_1\delta^3))}{H^{\delta + \frac{c_2\delta^5}{\ln\ln(H+2)}}}.$$

A completely similar result holds for any $\alpha \in \mathbb{Q}^*$, except that c_1 and c_2 now depend on α . Note that Mahler [10, p. 135, Satz 3] obtained a better exponent when $\alpha = 1$:

$$\left|\sum_{k=0}^{\delta} a_k e^k\right| \ge \frac{1}{H^{\delta + \frac{c_3 \delta^2 \ln(\delta+1)}{\ln \ln(H+2)}}}$$

where $c_3 > 0$ is absolute and $H \ge H(\delta)$ (which is not explicit).

b) We fix α of degree $d \geq 2$ over \mathbb{Q} . In this case, $p \approx 2d\delta$, $q \approx ce^{2d\delta}$, $t \approx c$, $a \approx e^{cd^2\delta^3}$, $b \approx e^{cd^2\delta^3}$, $u \approx e^{cd^2\delta^2}H^{2d}$, $v \approx e^{cd^2\delta^2}$ (where c denotes a constant that depends on α but not on d, and may differ in each of the previous estimates). It follows that there exist two constants $c_4, c_5 > 0$, that depend on α but not on d, such that for any $H \geq 1$, any integer $\delta \geq 1$ and any vector $(a_0, \ldots, a_\delta) \in \mathbb{Z}^{\delta+1} \setminus \{0\}$ with max $|a_k| \leq H$, we have

$$\left|\sum_{k=0}^{\delta} a_k e^{k\alpha}\right| \ge \frac{\exp(-\exp(c_4 d^2 \delta^2))}{H^{\psi(d,\delta,\lambda) + \frac{c_5 d^4 \delta^3}{\ln\ln(H+2)}}}.$$

c) If $\delta = 1$, then p = 2d-2 and $\psi(d, \delta, 2d-2) = 4d^2 - 2d - 1$. Estimating the parameters as above and using the fact that $d(1/\alpha)$ and $\overline{1/\alpha}$ are both bounded by $\sqrt{d+1}H(\alpha)$, we deduce the existence of two absolute constants $c_6, c_7 > 0$ such that for all $\alpha \in \overline{\mathbb{Q}}^*$ of height $H(\alpha)$ and degree d, for all $(p,q) \in \mathbb{Z} \times \mathbb{N}^*$, we have

$$\left|e^{\alpha} - \frac{p}{q}\right| \ge \frac{\exp\left(-\exp(c_6 s(\alpha)d)\right)}{q^{4d^2 - 2d + \frac{c_7 d^3 s(\alpha)}{\ln\ln(q+2)}}},$$

where $s(\alpha) := d + \ln H(\alpha)$ is the classical quantity called the size of α (see [3]). This is an explicit version of Kappe's irrationality measure [9].

4 A more general construction

We conclude this paper with a more general construction of K-linear approximations to powers of e^{α} . However, we shall prove that it does not lead to better results than those obtained in the previous sections.

We start as in §2, of which we borrow the notations. For any integer $1 \le k \le p - \delta + 1$, the k vectors ${}^{t}(a_0, \ldots, a_{\delta}, 0, \ldots, 0)$, ${}^{t}(0, a_0, \ldots, a_{\delta}, 0, \ldots, 0)$, $\ldots, {}^{t}(0, \ldots, 0, a_0, \ldots, a_{\delta}, 0, \ldots, 0)$ of \mathbb{C}^{p+1} are \mathbb{C} -linearly independent. We complete them with the p - k + 1 \mathbb{C} -linearly independent vectors ${}^{t}(A_{0,\ell}, A_{1,\ell}, \ldots, A_{p,\ell})$ ($\ell = 0, \ldots, p - k$) of \mathbb{C}^{p+1} to form a basis of \mathbb{C}^{p+1} . The previous construction corresponds to the case $k = p - \delta + 1$.

The integers p, d, δ, k must satisfy the conditions:

$$d \ge 1$$
, $p \ge \delta \ge 1$, $1 \le k \le p - \delta + 1$, $p + kd - d(p+1) + 1 > 0$.

We no longer require that $p \ge d\delta$.

Then similar "determinantal" computations as in the proof of Theorem 1 show that $|L_0| \ge c H^{-\varphi(d,\delta,p,k)-\varepsilon}$ where

$$\varphi(d, \delta, p, k) := dk - 1 + \frac{d^2k(p - k + 1)}{p + kd - d(p + 1) + 1}$$

The parameter δ does not appear explicitly in φ , and only through the above inequalities.

If d = 1, $\varphi(d, \delta, p, k) = p$ is minimal for $p = \delta$ (as expected) and it is enough to take k = 1 to obtain this result in this case.

If $d \geq 2$, we have

$$\frac{\partial \varphi}{\partial k} = -\frac{d(d-1)(p+1)^2}{(p+kd-d(p+1)+1)^2} < 0.$$

Therefore, for fixed $d \ge 2$ and $p, \delta \ge 1$, $\varphi(d, \delta, p, k)$ is minimal for $k = p - \delta + 1$ in which case it coincides with $\psi(d, \delta, p)$.

In both cases, this justifies to consider only the case $k = p - \delta + 1$ in the previous sections.

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